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Presence of inertial induction in general relativity

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Absence of Inertial Induction in General Relativity

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I review arguments indicating that there is no real, physically detectable, local inertial-induction effect in general relativity, contrary to recent comments by Tittle.

In a recent Letter Tittle¹ has brought up an old suggestion of Einstein's that there is some sort of inertial-induction effect in his standard general-relativistic theory of gravitation. In his book Einstein² devoted about ten pages to a discussion

of this problem. Dicke and I were not satisfied that general relativity met this criterion. In fact, we came to the conclusion that Einstein's claim of inertial induction was a purely coordinate effect and thus could have no physically detectable consequences.



What is inertia?

It takes work to accelerate a massive object. What is pushing back?

$$F = M_i \frac{d\mathbf{v}}{dt}$$

acceleration = (applied force) / (inertial mass)

Inertia involves a local momentum transfer!

Relativistic Gravitational Force Equation

– *geodesic equation* –

the separation of gravity and inertia: “a mischief of classical mechanics”

$$m_i \frac{dv}{dt} = m_g \frac{MG}{r^2}$$

Newton
2nd Law

Newton
Law of
Gravity

$$U^\nu \nabla_\nu U^\mu = \frac{dU^\mu}{d\tau} + \Gamma_{\alpha\beta}^\mu U^\alpha U^\beta = 0$$

tensor

“inertia”
not a tensor

“gravity”
not a tensor

$$U^\mu \equiv \frac{dx^\mu}{d\tau}$$

$$c^2 d\tau^2 \equiv g_{\mu\nu} dx^\mu dx^\nu$$

Linear gravity, metric components grouped by rotation properties

$$g_{\mu\nu} \simeq \eta_{\mu\nu} + h_{\mu\nu}, \quad |h_{\mu\nu}| \ll 1$$

$$h_{tt} \equiv -2\phi \rightarrow \text{scalar}$$

$$h_{ti} \equiv w^i = \mathbf{w} \rightarrow 3\text{-vector}$$

$$h_{ij} \rightarrow 3\text{-tensor}$$

Linearized geodesic equation

linear in metric perturbations – linear in test particle speed

$$\frac{1}{e} \frac{d\mathbf{p}}{dt} \simeq -\nabla\phi - \frac{1}{c} \partial_t \mathbf{w} + \mathbf{v} \times \nabla \times \mathbf{w} - \frac{v^n}{c^2} \partial_t h_{mn} \quad *$$

$$e = \gamma mc^2, \quad \mathbf{p} = \gamma m \mathbf{v}, \quad \gamma \simeq (1 - 2\phi + 2\mathbf{w} \cdot \mathbf{v}/c)^{-1/2}$$

gravito-electric and gravito-magnetic fields

$$\mathbf{E}_g \equiv -\nabla\phi - \partial_t \mathbf{w}/c, \quad \mathbf{B}_g \equiv \nabla \times \mathbf{w}$$

field equations

$$\nabla^2 \phi = \frac{\kappa}{2} \rho c^2$$

$$\nabla^2 \mathbf{w} = -2\kappa \rho c \mathbf{U}$$

Einstein's (1921) enumeration of Machian features of GR, based on the truncated linear force equation*

$$\frac{1}{c^2} \frac{d}{dt} [(1 + \phi) \mathbf{v}] = -\nabla \phi - \frac{1}{c} \partial_t \mathbf{w} + \mathbf{v} \times (\nabla \times \mathbf{w}) - \frac{v^n}{c^2} \partial_t h_{mn}$$

“1. The inert mass is proportional to $1+\Phi$, and therefore increases when ponderable masses approach the test body”.

“2. There is an inductive action upon accelerated masses, of the same sign, upon the test body”.

“3. A material particle, moving perpendicular to the axis of rotation inside a rotating hollow body, is deflected in the sense of the rotation...as shown by Thirring”.

*Einstein 1921,
Davidson 1957,
Brans 1962,
Brans 1977,
Ciufolini & Wheeler
1995

Brans on inertial induction (1962, 1977): Einstein 1

- “Absence of inertial induction in general relativity”, 1977
- **referred to effect No. 1** from Einstein 1921, an apparent effect on inertial mass from nearby matter, although Einstein used the word “induction” for effect No. 2.
- Never calculated per se, only argued that the Equivalence Principle prevents an external mass from influencing local inertia, because one can always transform to a frame in free fall, where the local mass vanishes.

Sciama (1954) Inertia from gravity (Einstein 2)

*weak-field
force
equation*

$$\frac{d\mathbf{p}}{dt} = mc^2 \left[-\nabla(\phi_U + \phi_L) - \frac{1}{c} \partial_t \mathbf{w} + \mathbf{v} \times (\nabla \times \mathbf{w}) \right] \rightarrow \text{set} = 0$$

1. evaluated in frame of accelerated object
2. total grav force from universe is zero

*cosmic
grav.
potential*

$$\phi_U = \int_0^{R_U} G \frac{\rho_U}{r} 4\pi r^2 dr = 2\pi G \rho_U R_U^2 = 2\pi G \rho_U c^2 / H^2 \sim c^2$$

$$H^2 = \frac{8\pi G}{3} \rho_U$$

Hubble constant H

*cosmic
vector
potential*

$$\mathbf{w} = \int_0^{R_U} \mathbf{v} \frac{\rho_U}{r} 4\pi r^2 dr = \phi_U \mathbf{v}$$

$$-c^2 \nabla \phi_L = \partial_t \mathbf{v}$$

“Newton Law of Gravity”

Inertia arises from gravitational induction;
and if so, then it predicts the mass of the universe

Linear gravity cannot describe inertia

$$h_{\mu\nu} \sim \Phi_U \sim 1$$

- Sciama inertial induction shown only for vector gravity
- Vector gravity only approximates tensor gravity for weak fields
- The gravitational field of the universe is large

Sciama-style inertial induction not proven in GR

C&W on inertial induction: Einstein 3

- Ciufolini & Wheeler (1995) – Gravitation and Inertia – “*inertia here from mass there*”
 - attempted to treat inertia in GR, building on results in MTW. Settled on a gravitomagnetic interaction in the Kerr metric: Einstein 3
 - struggled with boundary conditions. Proposed closed universe with no cosmological constant, yet universe now understood to be 70% CC

$$ds^2 = -\left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 + \frac{2GM}{r}\right) dr^2 + r^2 d\Omega^2 - \frac{4J}{r} \sin^2 \theta d\phi dt$$

gravitomagnetism

Does inertia from gravity violate the EP?

~the main conceptual obstacle~

- Sciama showed inertia only for a vector theory of gravity
- Tensor gravity obeys the Equivalence Principle
 - all gravitational fields vanish locally in free fall
- How does a body couple to the gravitational field of the universe if it vanishes in free fall?

Local gravitational phenomena

– *Is it possible?* –

- If inertia is local, and if it is gravitational, then it is a local gravitational phenomenon
 - such phenomena are unknown to GR, and at odds with the Equiv. Princ.
- If inertia is a local gravitational phenomenon arising from the grav field of the universe, then perhaps momentum can be exchanged locally with the grav field of the universe

Gravitational field of the universe: ~2005

– *flat Robertson-Walker metric* –

$$g_{\mu\nu}^U dx^\mu dx^\nu = c^2 dt^2 - a^2(t)[dx^2 + dy^2 + dz^2]$$

$$g_{\mu\nu} = \begin{bmatrix} c^2 & 0 & 0 & 0 \\ 0 & -a^2 & 0 & 0 \\ 0 & 0 & -a^2 & 0 \\ 0 & 0 & 0 & -a^2 \end{bmatrix} \quad H^2 \equiv \left(\frac{da/dt}{a} \right)^2 = \frac{8\pi}{3} \rho_U G \equiv t_U^{-2}$$

$$\Phi_U = \int_0^{ct_U} \frac{\rho_U G}{r} 4\pi r^2 dr \propto \rho_U G c^2 t_U^2 \sim c^2$$

Hubble Drag

~force from cosmic metric on a body in motion~

$$U^\nu \nabla_\nu U^\mu = \frac{dU^\mu}{d\tau} + \Gamma_{\alpha\beta}^\mu U^\alpha U^\beta = 0$$

$$g_{\mu\nu}^U dx^\mu dx^\nu = c^2 dt^2 - a^2(t)[dx^2 + dy^2 + dz^2]$$

body with velocity v along
x-direction

$$U^\mu = (U^t, U^x, 0, 0) = (U^t, v U^t, 0, 0)$$

a body in motion
experiences drag force from
the Hubble expansion

$$\frac{dU^x}{dt} = -2 \frac{\dot{a}}{a} U^x = -2 H U^x$$

$$H_0 t \sim H_0 * 1000 \text{ s} = 10^{-14}$$

Force already known from galaxy dynamics.
Also applies to small bodies at a point.

Lorentz transformation (boosts) of metrics

~gravitational field in moving frame~

$$\tilde{g}_{\mu\nu} = g_{\alpha\beta} \Lambda_{\mu}^{\alpha} \Lambda_{\nu}^{\beta}$$

$$\Lambda_t^x = \Lambda_x^t = \gamma \frac{v}{c} , \quad \Lambda_t^y = \Lambda_y^t = \Lambda_t^z = \Lambda_z^t = 0$$

$$\Lambda_t^t = \gamma , \quad \Lambda_x^x = \gamma , \quad \Lambda_y^y = \Lambda_z^z = 1$$

$$\Lambda_j^i = 0 , \quad i \neq j$$

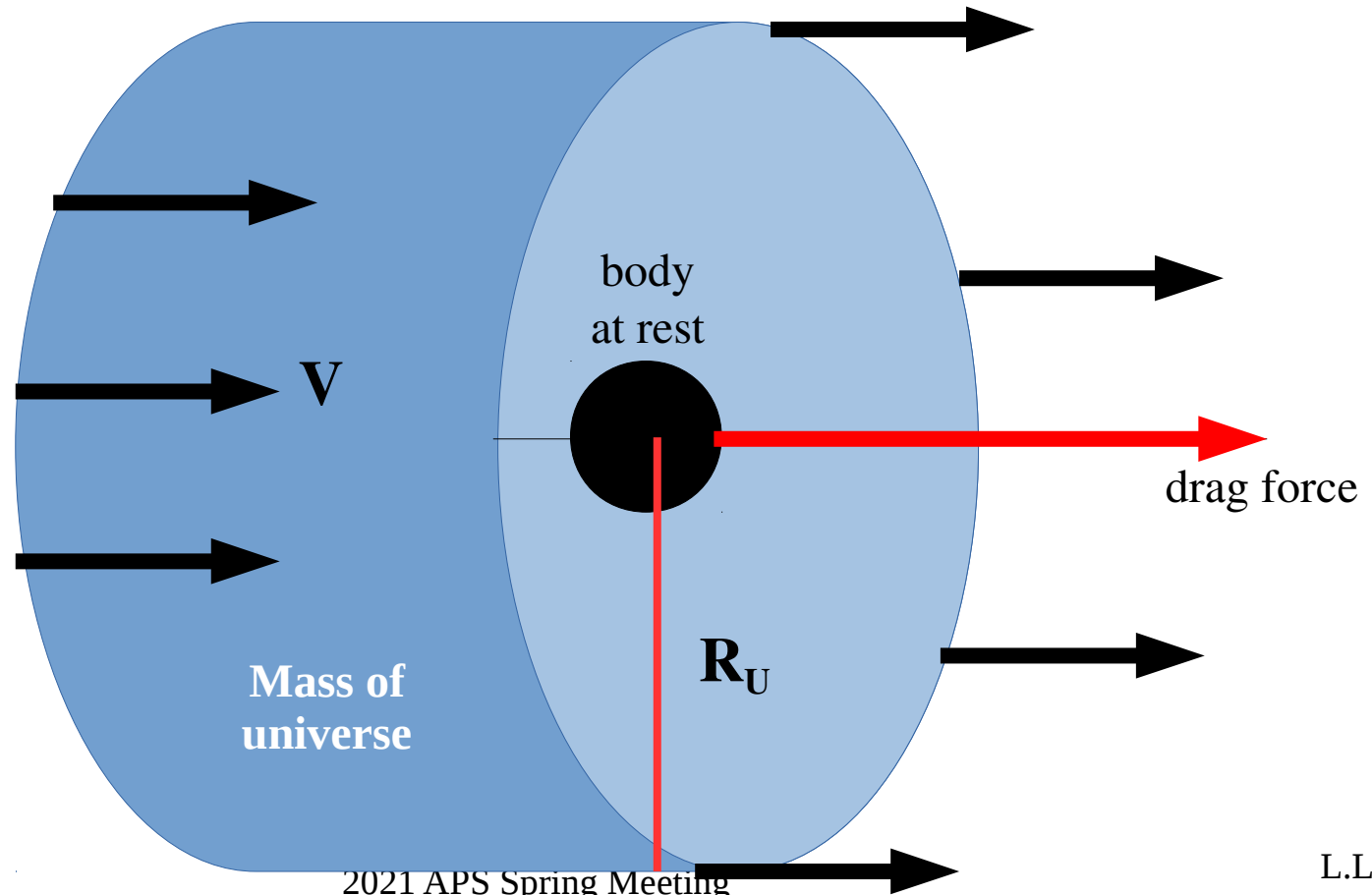
$$\tilde{\eta}_{\mu\nu} = \eta_{\alpha\beta} \Lambda_{\mu}^{\alpha} \Lambda_{\nu}^{\beta} = \eta_{\mu\nu}$$

$$g_{tx} = 0 , \quad \tilde{g}_{tx} \neq 0$$

Gravitomagnetism emerges in the boosted frame
of non-Minkowski metrics

Hubble drag concept from the boosted frame

- A moving object sees a surrounding mass current from the universe
- The motion creates a gravitomagnetic force in the boosted frame
- Confirms Sciama's basic idea
- aka linear frame-dragging



Identification of Rectilinear Frame Dragging

- the Hubble drag effect “looks” gravitomagnetic in all cosmic reference frames except the isotropic frame
- the rest-frame force on moving bodies can be considered frame dragging
- this is the first specific example of rectilinear frame dragging in the literature. Classic frame dragging involves rotating mass fluxes.

Conclusions

- General relativity implies at least 3 non-Newtonian effects on inertia
- They have been conflated in the literature. Brans showed that one was a coordinate effect, but did not consider the other two effects
- The known gravitational field of the universe, the Robertson-Walker metric, induces gravitomagnetic components in boosted frames
- These gravitomagnetic effects lead to a drag force known from galaxy dynamics
- This is the first reported case of linear-geometry frame dragging. Previous work has been for rotational geometries.
- This implies the cosmic gravitational field can be detected in principle locally through inertial-frame experiments with small bodies.